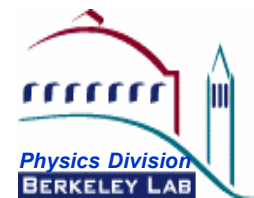


Cosmology Teach-In / July 3, 13 700 002 003 ABB

Polarization of the Cosmic Microwave Background Radiation

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Outline

- A photon's journey through the universe
- The CMB Power Spectrum
- Polarization at the surface of last scattering
- Primordial Polarization from Inflationary Gravitational Wave Background
 - ✿ secondaries: grav. lensing
- Experiments for polarization
 - ✿ road to measuring gravity wave signal

Our Habitat

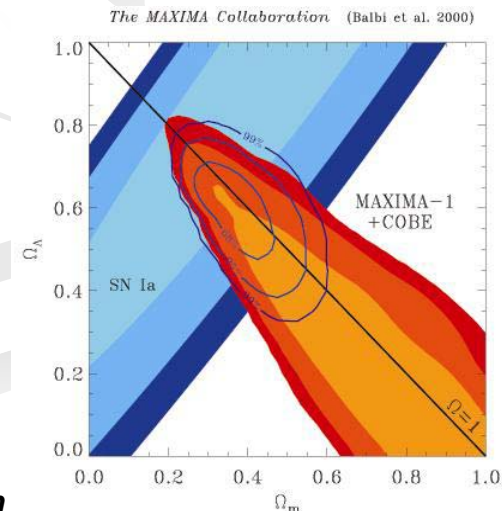
“my first reaction to cosmology was one of surprise that grown people could seriously care about such a schematic physical theory.” (Peebles)

CMB temperature anisotropy measurements of the acoustic peaks (+SN1a) have revolutionized our self image

speculative cosmology has evolved into our cosmological model

- we live in a spatially flat (critically dense) $\Omega_{\text{tot}}=1$ universe
- consisting mainly (95%) of dark matter and dark energy
- in which structure arises through *gravitational instability* from quantum fluctuations during an inflationary epoch.

a “*Cosmic Concordance*” has emerged.



The Electromagnetic Spectrum

cold

Long Wavelength
Low Frequency
Low Energy



Aircraft and
Shipping
Bands



AM
Radio

Shortwave
Radio



TV and
FM Radio



Microwaves
Radar



Infrared
Light

cm \rightarrow km
 $10^{-2} - 10^3$ m

mm \rightarrow cm
 $10^{-3} - 10^{-2}$ m

μ m
 $10^{-6} - 10^{-4}$ m

100s nm
 $10^{-7} - 10^{-6}$ m

$1 - 10^2$ eV

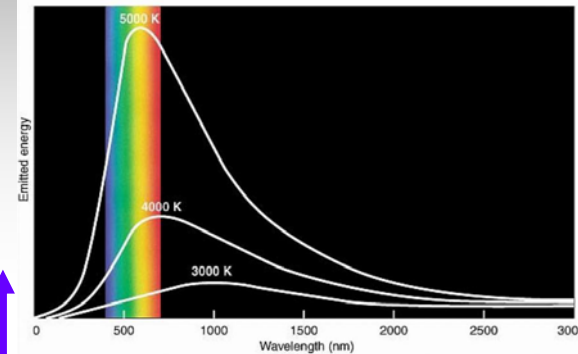
$10^2 - 10^5$ eV

$> 10^5$ eV

"CMB" photons
in the future.

CMB
Photons
Now

"CMB" photons
a long time ago.



as the universe
expands & cools, the
Cosmic Background
Radiation photons
are stretched,
and evolve through
the EM
spectrum.

hot

Short Wavelength
High Frequency
High Energy



Ultraviolet
Light



X-rays



Gamma-rays

Life of a Cosmic Photon

Inflation

- at $t \sim 10^{-38}$ s ABB, the universe undergoes a phase transition causing an explosive e^{60} exponential expansion
 - leaves its imprint as *inflationary gravitational wave background* (IGB)
 - to be **encoded on the CMB** at the surface of last scattering as a curl (***B-mode***) component to the polarization
- inflation provides answers for:
- horizon problem: explains why CMB is isotropic on cosmically unconnected scales $\geq 1-2^\circ$
 - flatness problem ($\Omega=1$): exponential expansion locally flattens spatial curvature to high precision.
 - predicts scale invariant adiabatic density perturbations

Timeline (ABB)

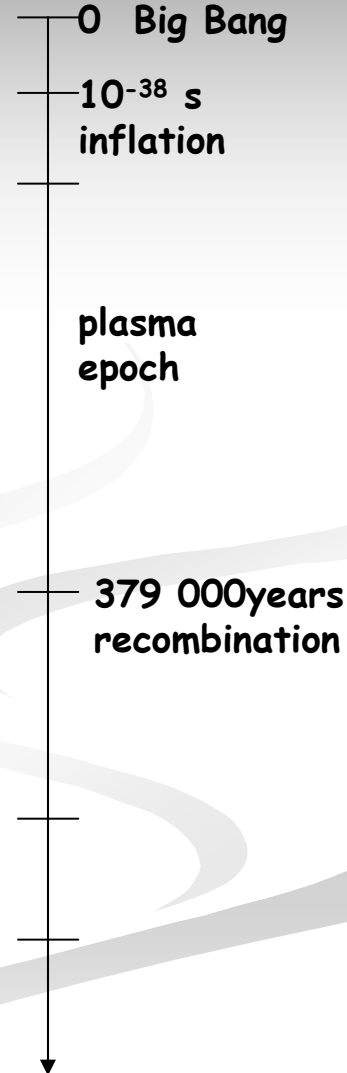


Life of a Cosmic Photon

Plasma Epoch

- density perturbations undergo gravitationally driven *acoustic oscillations*.
 - ✿ gravitational wells from matter overdensities cause the plasma to fall inwards
 - ✿ until eventually it rebounds under radiation pressure.
- matter and radiation are *tightly coupled*
 - ✿ information about matter distributions are precisely encoded in radiation

Timeline (ABB)



Life of a Cosmic Photon

$t < \sim 400\,000$ years

- matter and radiation are *tightly coupled*
 - ⊗ electrons are Coulomb interaction coupled to baryons
 - ⊗ photons *Thomson scatter* off electrons
 - mean free path $\lambda_c = 2.5$ Mpc (short!)
 - the universe is *opaque* on cosmological scales.

Recombination $t \sim 379\,000$ years ABB $z \sim 1089 \pm 1$ $T = 3000$ K

- protons combined with electrons to form neutral hydrogen
 - ⊗ recombination is sudden, $\Delta z \sim 195 \pm 2$.
- this is the “surface of last scattering” for most photons- which travel unhindered to us, providing a snap-shot of the universe at time = 379 K years.

Timeline (ABB)

0 Big Bang

10^{-38} s
inflation

plasma
epoch

379 000 years
recombination



Life of a Cosmic Photon

Large Scale Structure Formation Epoch

- Matter collapses under gravity to form the rich structure (including us!) of the universe.

Timeline (ABB)

— 0 Big Bang

— 10^{-38} s
inflation

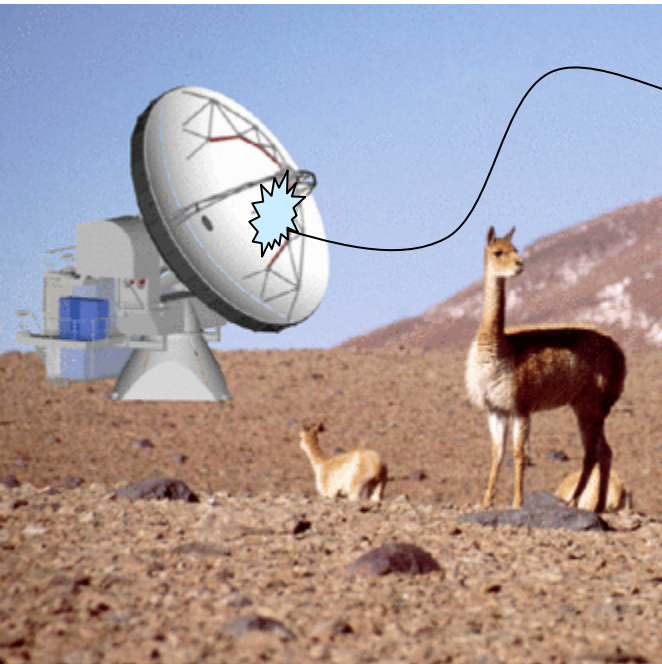
- Why the difference between Matter & Radiation now??
 - ⊗ matter and radiation had the same distribution a few $\cdot 10^5$ years ABB, why not now?
 - *radiation pressure* resists the pull of gravity.
 - ⊗ this is calculationally important!, because the perturbations in the uniform photon distribution are small, ***linear perturbation theory applies*** and we can calculate things.
 - ⊗ CMB observations are important not just because they carry a lot of information... but also because ***they are theoretically easy (robust) to predict.***
(analogy between EW & QCD).

Life of a Cosmic Photon

NOW

✿ $t \sim 13\,700\,002\,003$ years ABB

✿ $T = 2.7$ K



End of the line for a precious few of the photons, as we trap them in our detectors and glean a little information about the universe we live in.

Timeline (ABB)

0 Big Bang

10^{-35} s
inflation

plasma
epoch

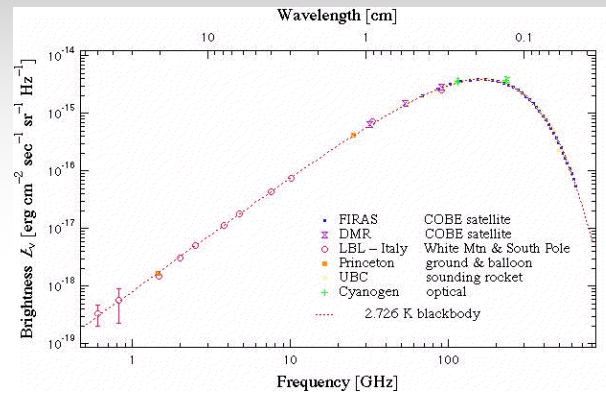
400 000yr
recombination

LSS
epoch

180 Myears
 $z \sim 20$
Reionization

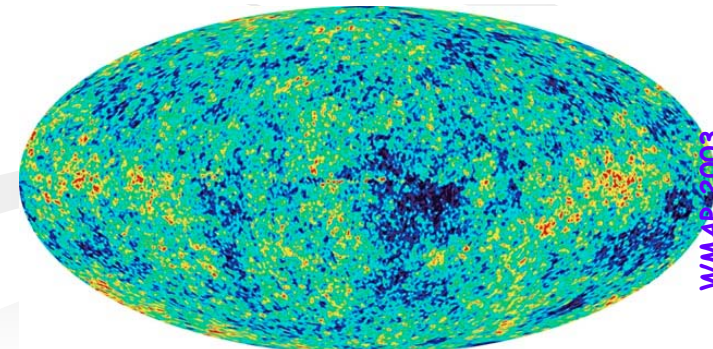
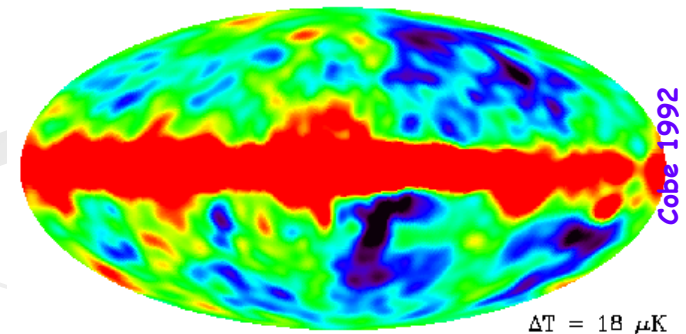
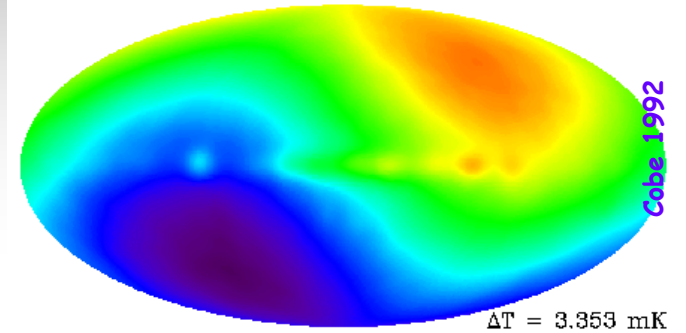
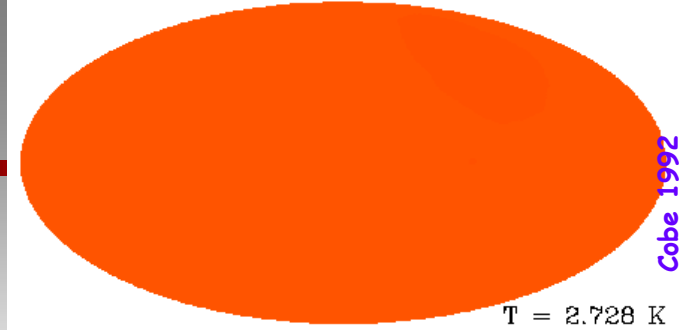
14.5 Gyears
NOW

- CMB is a near perfect black body, 2.7° K ($\neq 0$ monopole)

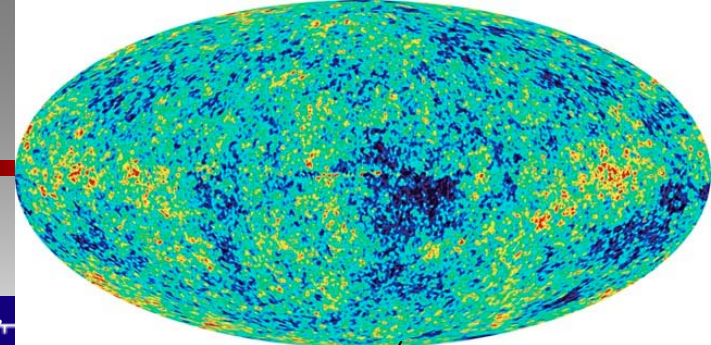


- Dipole Anisotropy 10^{-3} ($\neq 1$ dipole)
(our motion in the CMB 'rest frame')

- Temperature Anisotropy 10^{-5}



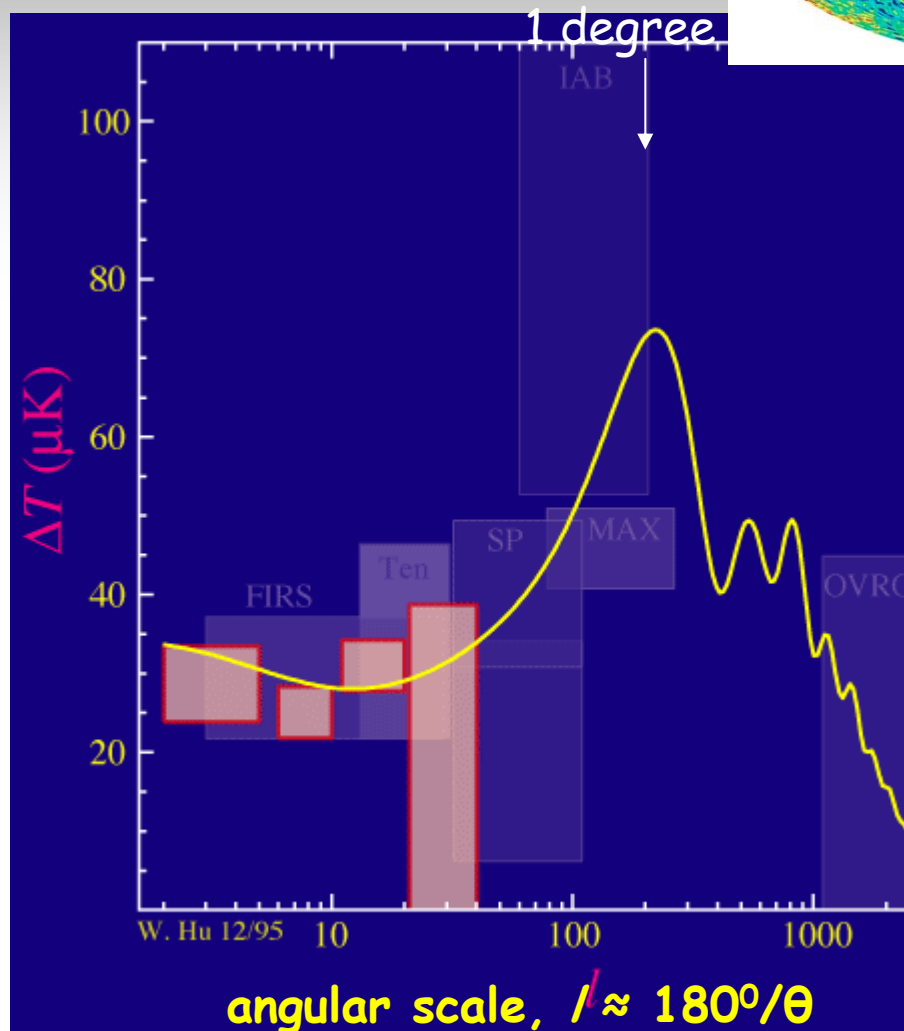
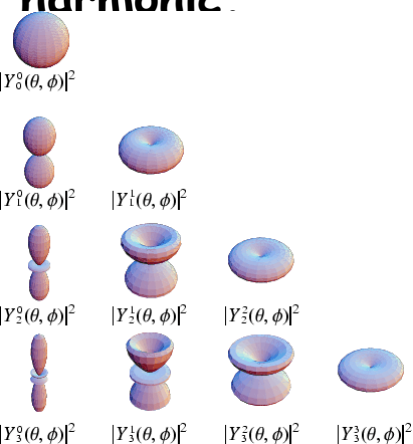
The Power Spectrum



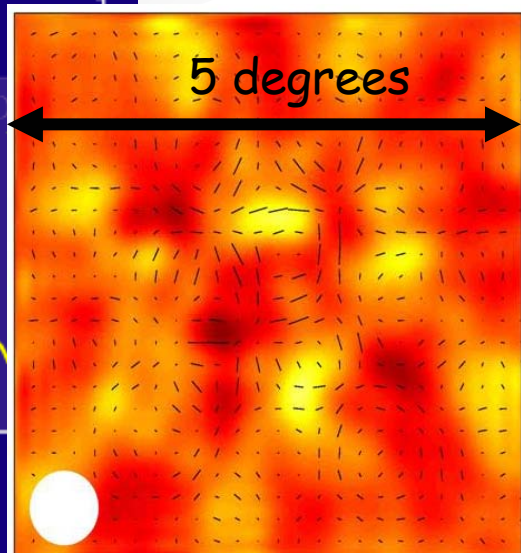
blotchy-ness
of the sky at a
particular scale

or...

how much the sky
looks like a
particular spherical
harmonic.

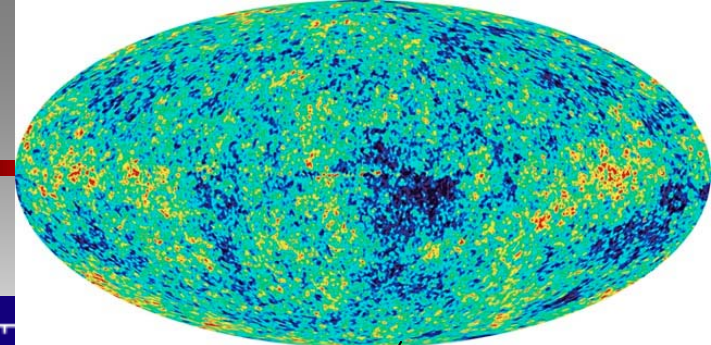


Gaussian fluctuations
→ all stat. info is
contained in the C_l 's.

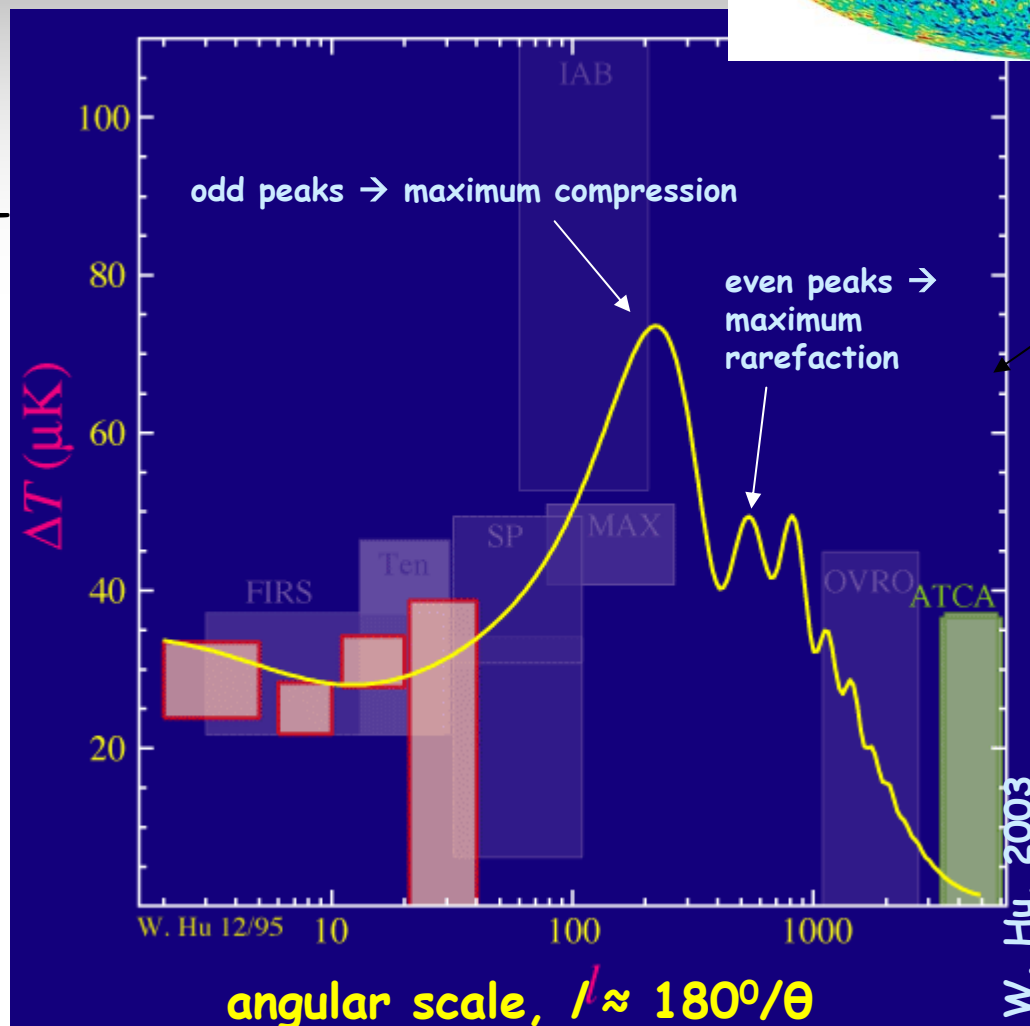


smaller scales

The Power Spectrum



Harmonic peaks
→ extrema in photon-
baryon fluid
oscillation cycle.



smaller scales

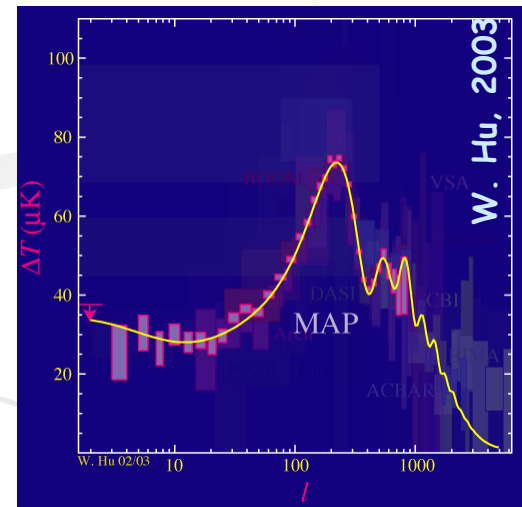
Measuring the CMB Temperature

- basic observable: intensity as a function of frequency and position
- for the temperature field,
 - ⊗ (assume Gaussian fluctuations)

$$\Theta(\hat{n}) = \frac{\Delta T}{T} \quad \text{temperature fluctuations}$$

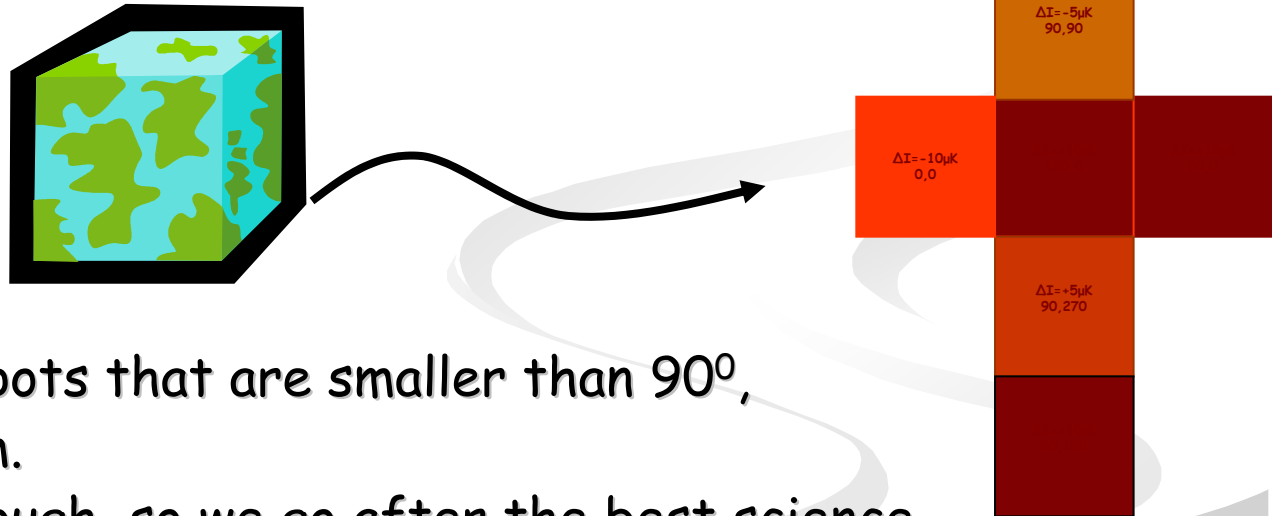
$$\Theta_{lm} = \int d\hat{n} Y_{lm}^*(\hat{n}) \Theta(\hat{n}) \quad \text{multipole moments}$$

$$\langle \Theta_{lm}^* \Theta_{l'm'} \rangle = \delta_{ll'} \delta_{mm'} C_l \quad \text{power spectrum}$$

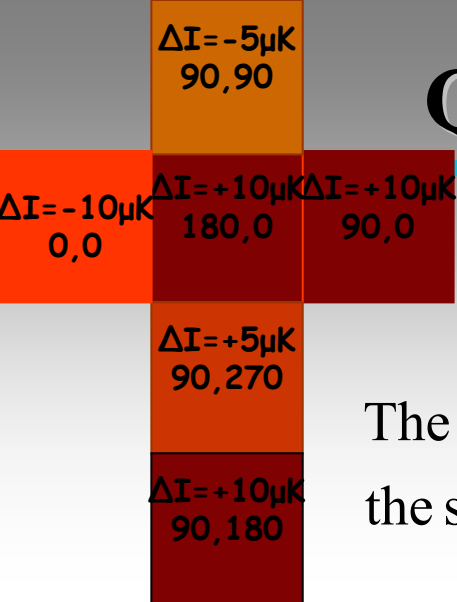


- ⊗ neglecting sky curvature, this reduces to a 2d-Fourier analysis.

- Consider a toy experiment with 90° beams (watch glass primary)
- no polarization sensitivity (yet)
- full sky coverage.
- We divide the sky into a cubic surface with 6 pixels, which we look out on from the inside:



- if the universe has spots that are smaller than 90° , we can't resolve them.
- But hey!, times are tough, so we go after the best science we can, and try to observe the lower l moments of the T anisotropy.



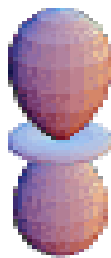
Quadrupole component of Power Spectrum

Our pixels are:
 $n=(\theta, \Phi) = 0, 0$
 $= 90, 0$
 $= 90, 90$
 $= 90, 180$
 $= 90, 270$
 $= 180, 0$

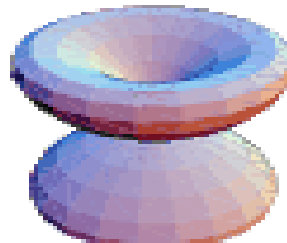
$$\Theta(\hat{n}) = \frac{\Delta T}{T}$$

The (complex) multipole moments quantify the extent to which the sky looks like the spherical harmonics :

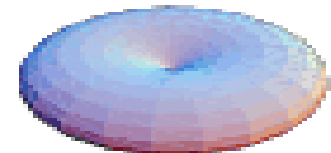
$$\Theta_{lm} = \int d\hat{n} Y_{lm}^*(\hat{n}) \Theta(\hat{n}) \stackrel{l=2, m=0}{=} \sum_{i=\text{pixel}} \frac{1}{2} \sqrt{\frac{5}{4\pi}} (3 \cos^2 \theta_i - 1) \cdot \frac{\Delta T_i}{T}$$



$$|Y_2^0(\theta, \phi)|^2$$



$$|Y_2^1(\theta, \phi)|^2$$



which we normally report in

$$\sqrt{\frac{l(l+1)C_l T^2}{2\pi}} = 62 \mu K.$$

Why this funny normalization?

Power per log interval in l

$C_l \propto 2l+1$ (since $m = -l, \dots, 0, \dots, +l$)

and the variance

$$= \int d^2 l \frac{C_l}{(2\pi)^2} = \int dl \frac{c \cdot 2l+1}{2\pi} = \frac{c \cdot l^2 + l}{2\pi} = \frac{l(l+1)}{2\pi} C_l$$

we've made some simplifications...



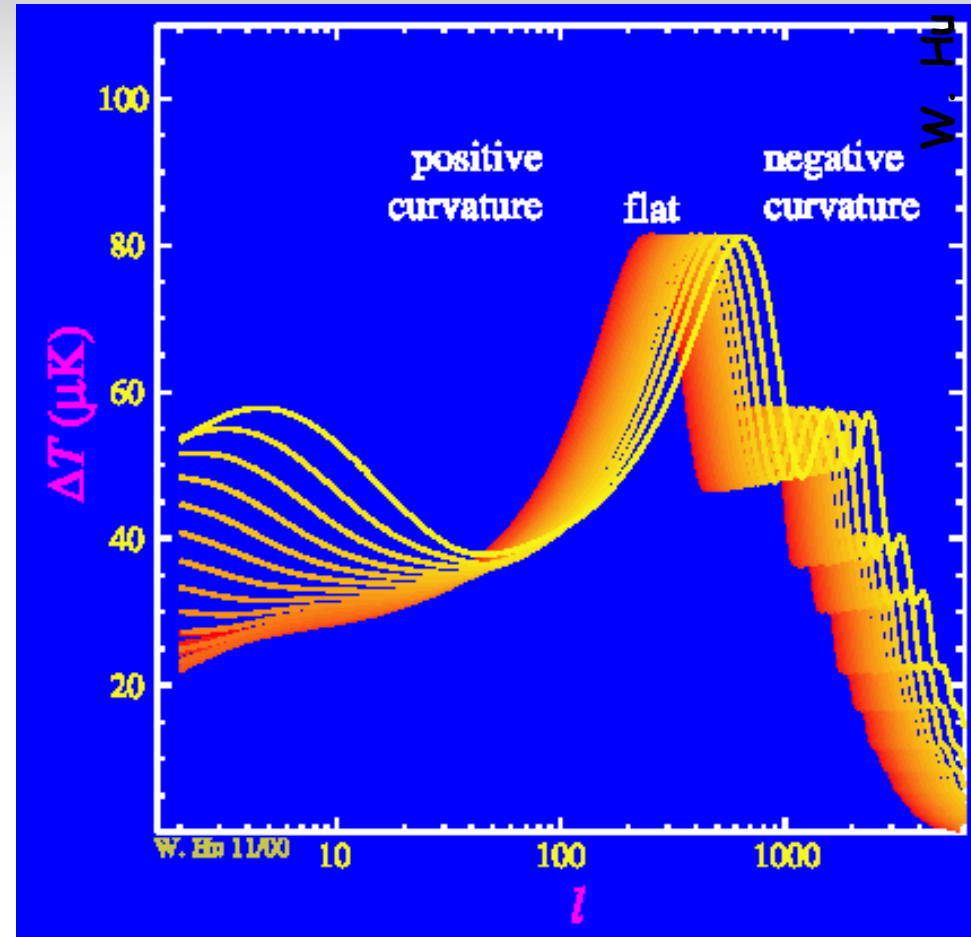
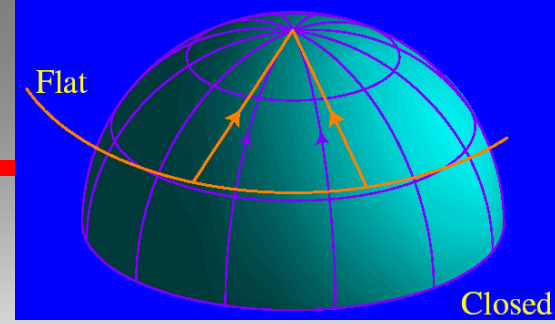
A field of physics cows

©2001 Drew Ness

- would need to de-convolute with beam shape functions, filter functions.
 - measure and account for noise
 - typically filter some portion of it out of the timestream
 - ⊗ correlated noise is a big pain.
- deriving the answer isn't difficult,.. but the computations typically go as $(\text{number of pixels})^3$, leading to computational problems/approximations

Curvature of Spacetime

- The angular location of the first peak measures the curvature of spacetime, $\Omega_{\text{tot}}=1$.
 - ✿ effectively measures if anisotropies travel straight out to us along curved paths.

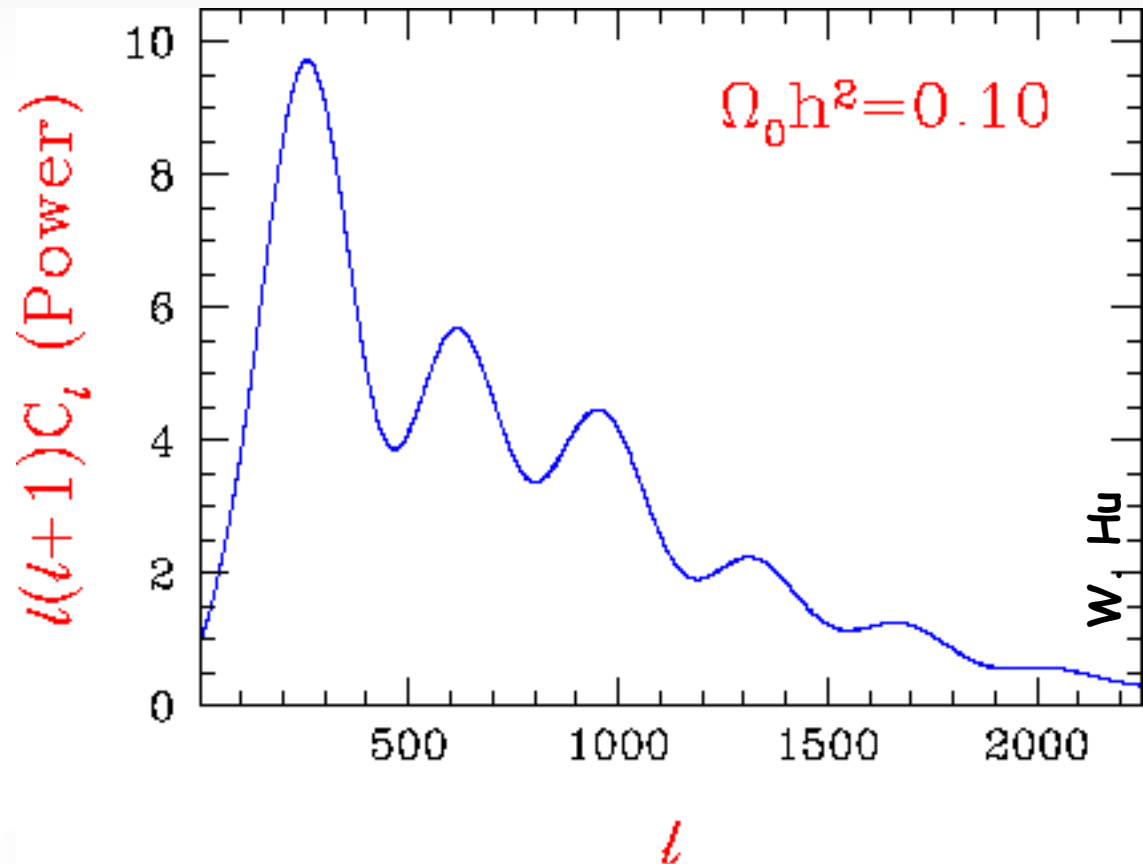


Power Spectrum: Ω_M

- total matter density

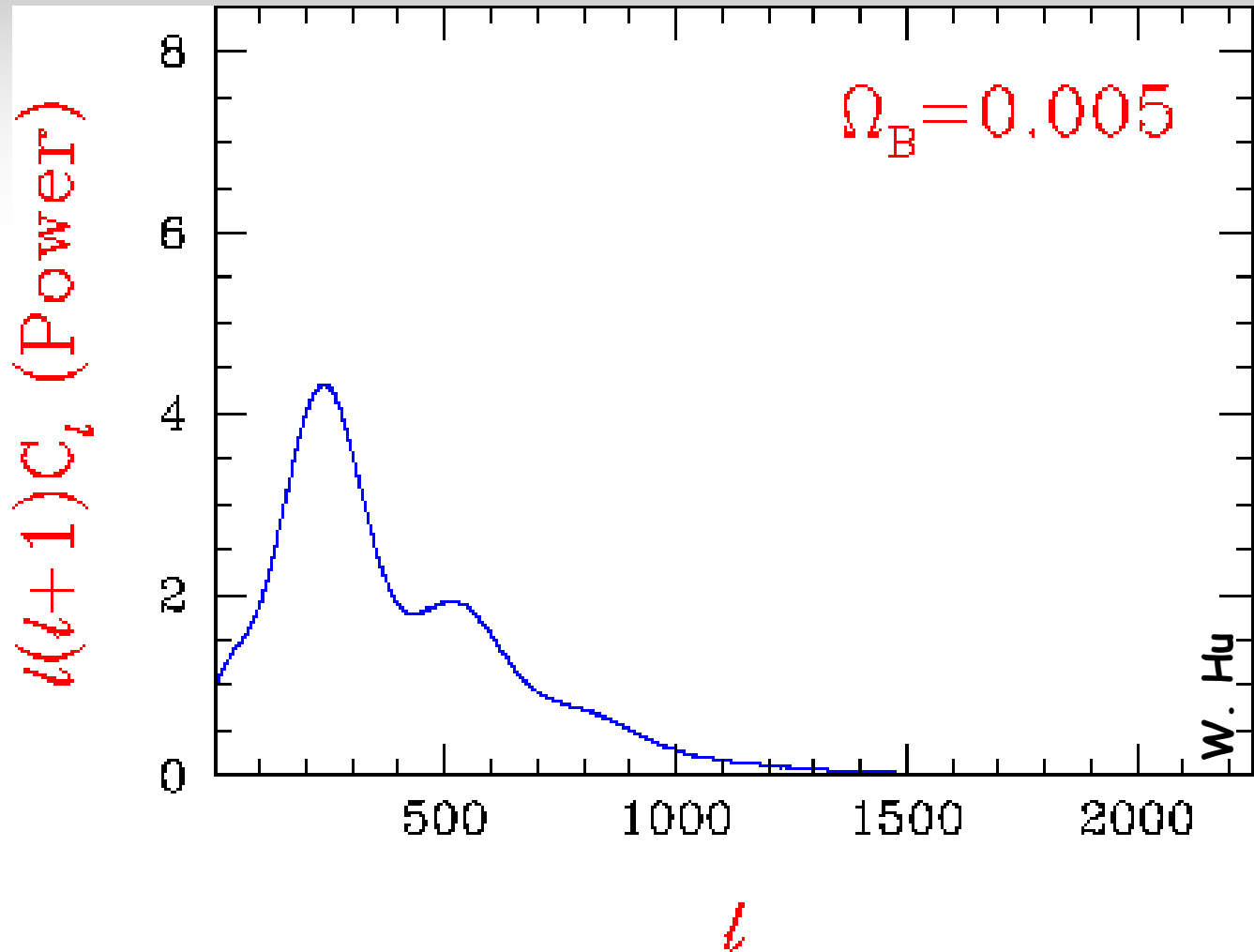
$$\Omega_M = \Omega_b + \Omega_{\text{CDM}}$$

determines the normalization of the power spectrum.

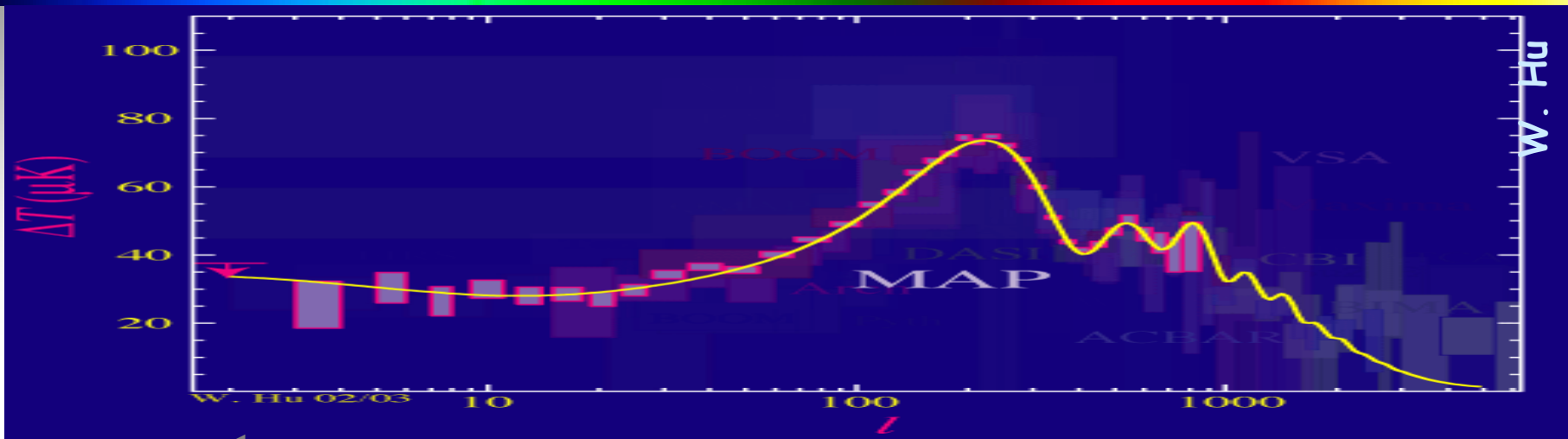


Power Spectrum: Ω_b

- ratio of 1st to 2nd peaks tells us the baryon content of the universe.
- enhances compression (odd) harmonic peaks.



What drives the errors?



Low l fundamental limitation is set by **Cosmic Variance**.

There is only one sky and only $2l + 1$ samples from the sky in each multipole moment.

Leads to an inevitable error

(further improvement by binning, $\Delta l \sim l$)

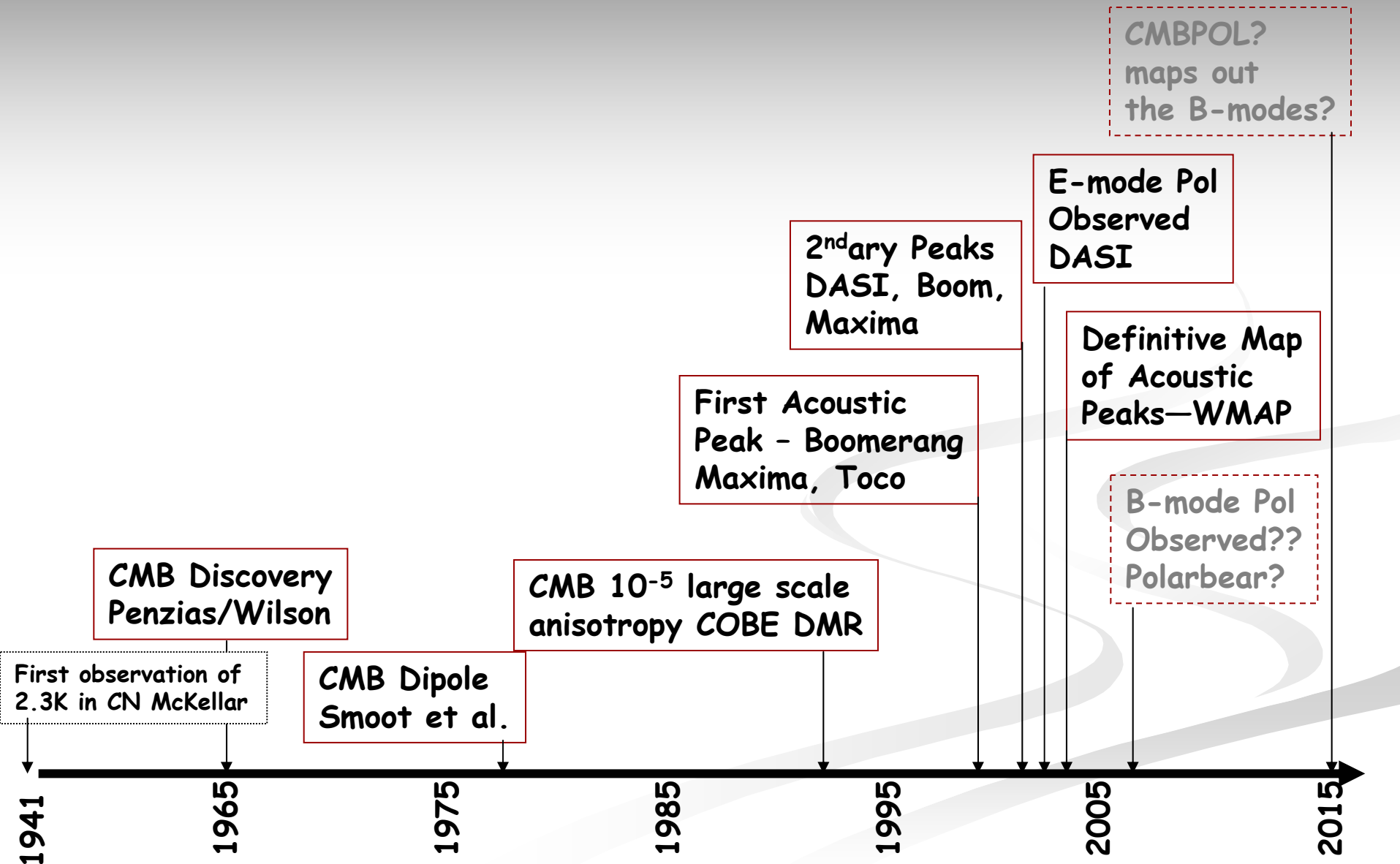
Example: WMAP is cosmic variance limited for $l < 354$.

$$\Delta C_l = \sqrt{\frac{2}{2l+1}} C_l$$

If we survey a fraction f_{sky} less than the full sky, **sample variance** increases the errors by $1/\sqrt{f_{\text{sky}}}$. (Planck $f_{\text{sky}} = .65$)

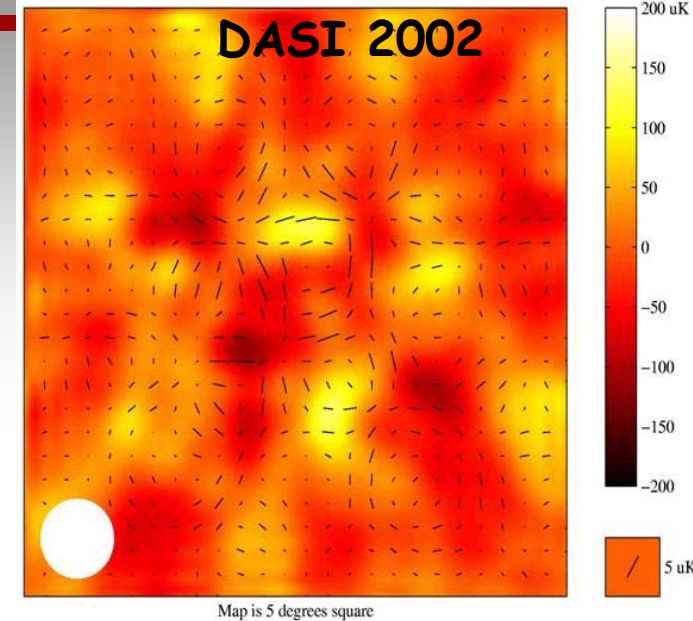
At high l , approaching the experiment resolution scale, noise (photon stats, electronics, astrophysical) dominate.

CMB Observations



Why measure CMB Polarization?

- scalar, vector & tensor fields carry more information than the temperature anisotropies alone.
- gives us more information about the acoustic peaks
- measure cosmo parameters better
- measure the reionization epoch, which produces a large degeneracy in the Temp spectrum
- measure gravity wave amplitude... the smoking gun of inflationary models.



Thomson Scattering & Polarization

- Incoming polarized light emerges polarized.

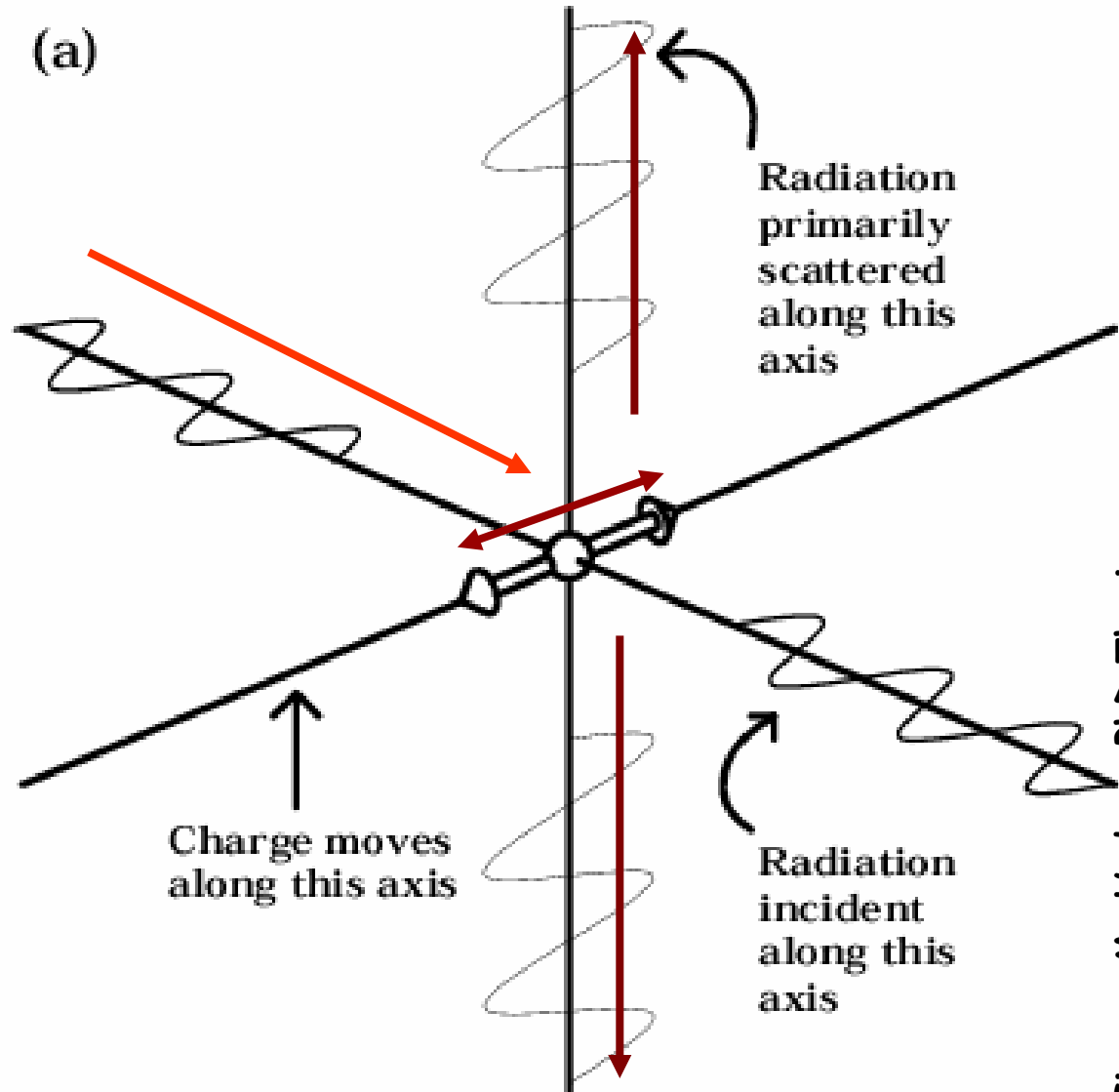
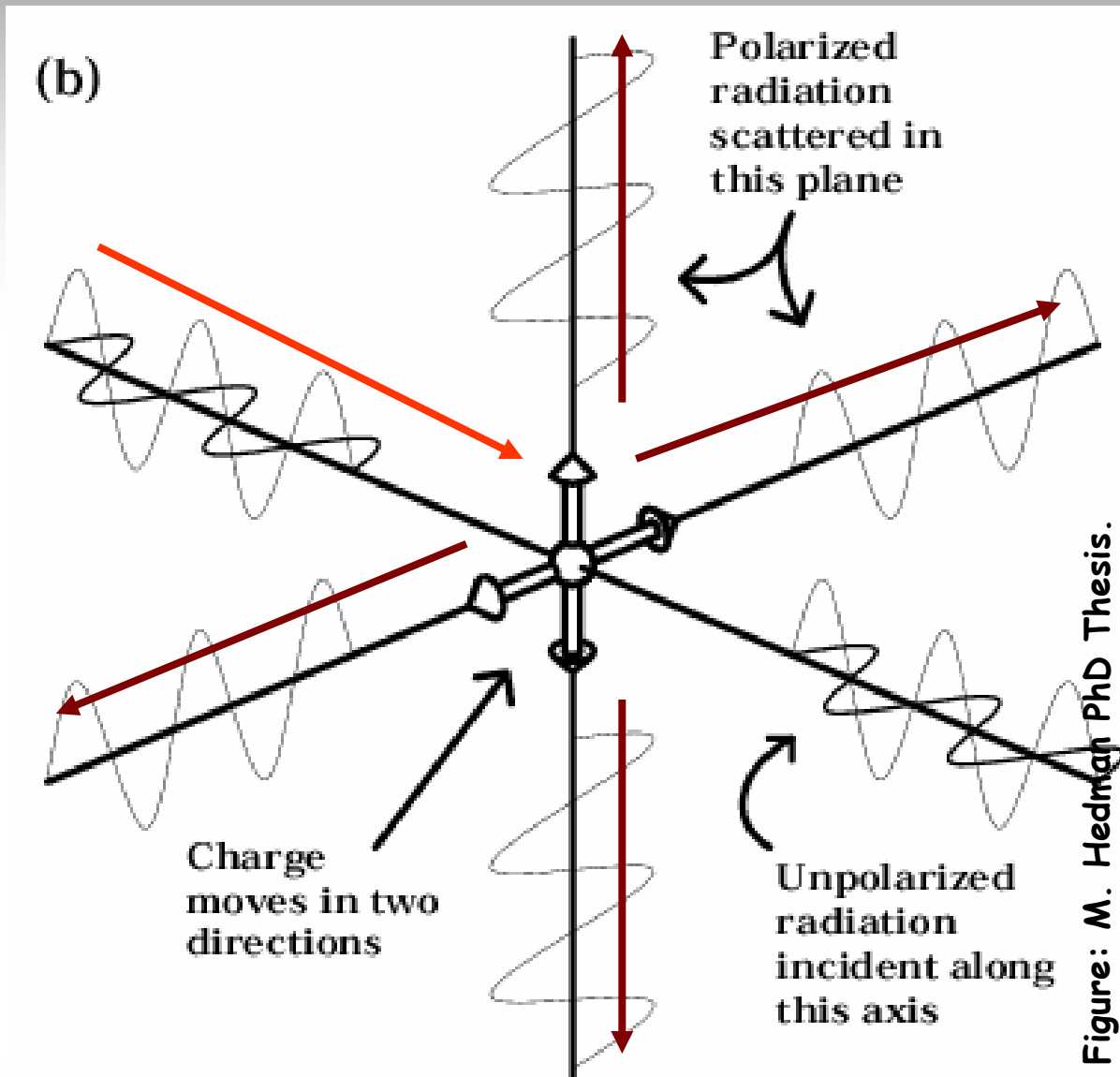


Figure: M. Hedman PhD Thesis.

Thomson Scattering & Polarization

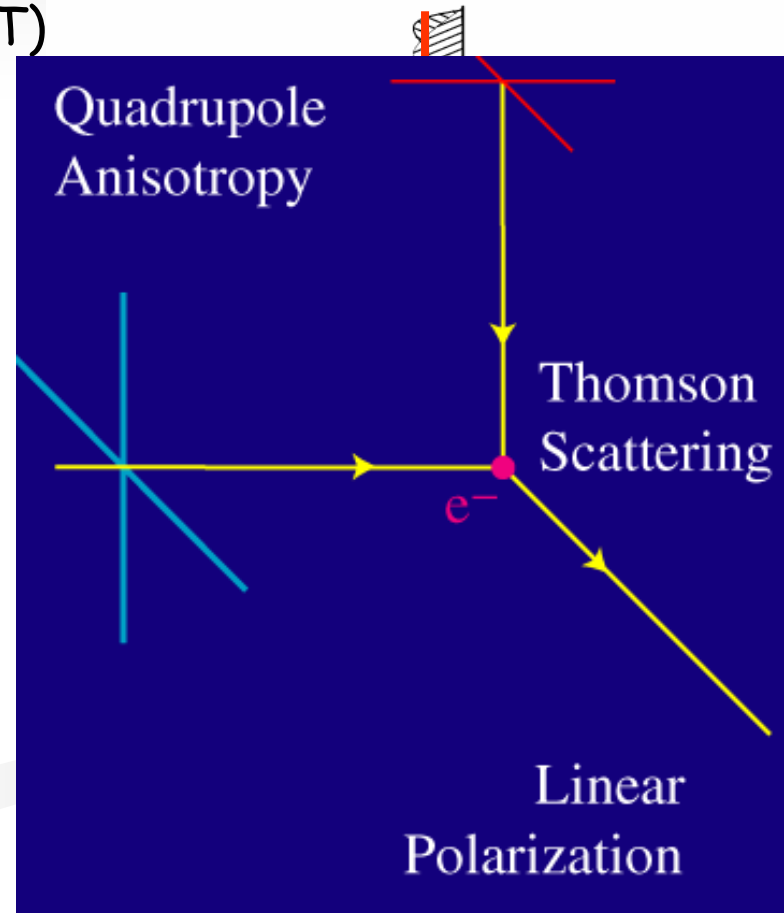
- A plane wave undergoing Thomson scattering produces polarized light too.
- but if we have equal amounts of light coming from all directions, there is no net polarization.



Thomson Scattering & Polarization

- Incoming polarized light emerges polarized.
- A plane wave undergoing Thomson scattering produces polarized light too.
- but if we have equal amounts of light coming from all directions, there is no net polarization.
- but if different intensities of quadrupole (unpolarized) light arrive from different directions, the net result is polarization. (10% of ΔT)

Figure: W. Hu doctored by G. Smoot



Polarized Light & Stokes Parameters

- Stokes vector $S=(I,Q,U,V)$ describes the intensity and polarization of light.

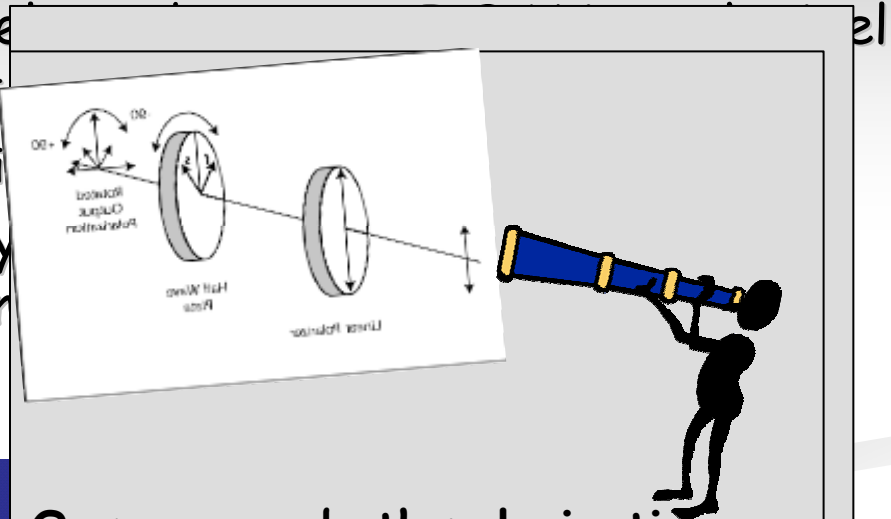
$$\begin{aligned} V &\equiv I_{\text{RCP}} - I_{\text{LCP}} \\ U &\equiv I_{+45} - I_{-45} \\ Q &\equiv I_0 - I_{90} \\ I &\equiv \text{total intensity} \end{aligned}$$

~~$V \equiv I_{\text{RCP}} - I_{\text{LCP}}$~~ absent in cosmo setting

- Think of Q as the “preference” of the light for horizontally ($Q=+1$) polarized light.
- For unpolarized light Q,U,V are all zero. For polarized light, $Q^2+U^2+V^2=1$.

Polarized Light

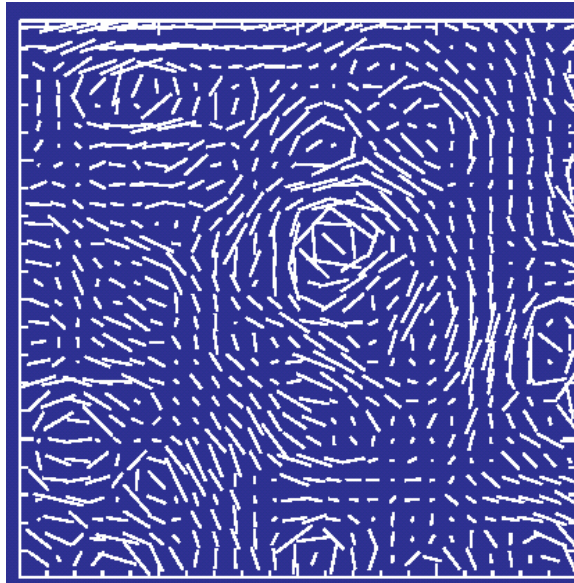
- Basic observable: intensity and polarization as a function of ν and n , $S=(I,Q,U,V)=S(\nu,n)$
- We can divide the sky up into pixels
 - ⊗ circular polarization, V is absent
 - ⊗ there is no E,B-mode info for individual pixels
- By doing a vector analysis of many pixels we can get the gradient ("E-mode") and curl ("B-mode")



Or measure both polarizations with the same detectors using a $\frac{1}{2}$ wave plate.

- removes many systematics
- but! throws away $\frac{1}{2}$ the light.

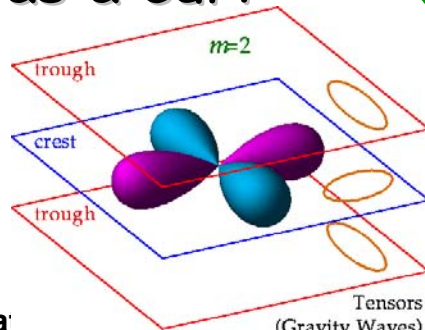
Separately measure E_x and E_y .



vector field
E (grad) mode
polarization

B (curl)-mode Polarization

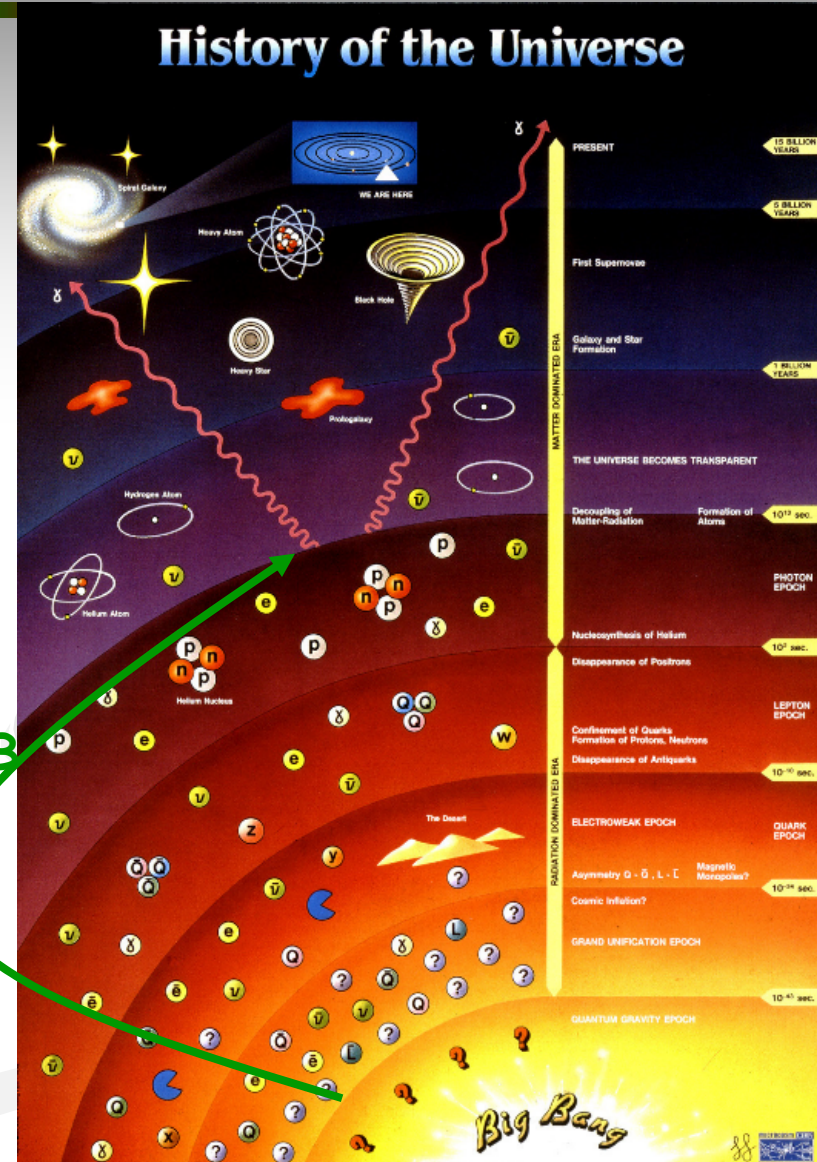
- Just as the CMB gives the fingerprint of the universe at $t=400$ Kyears,
- the fingerprint at the time of inflation (10^{-38} s) is encoded on the *inflationary gravity wave background* (IGB).
- This gravity wave fingerprint stretches and compresses spacetime on the surface of last scattering—transmitting that information as a curl component in the polarization.



Ma

Tensors
(Gravity Waves) >

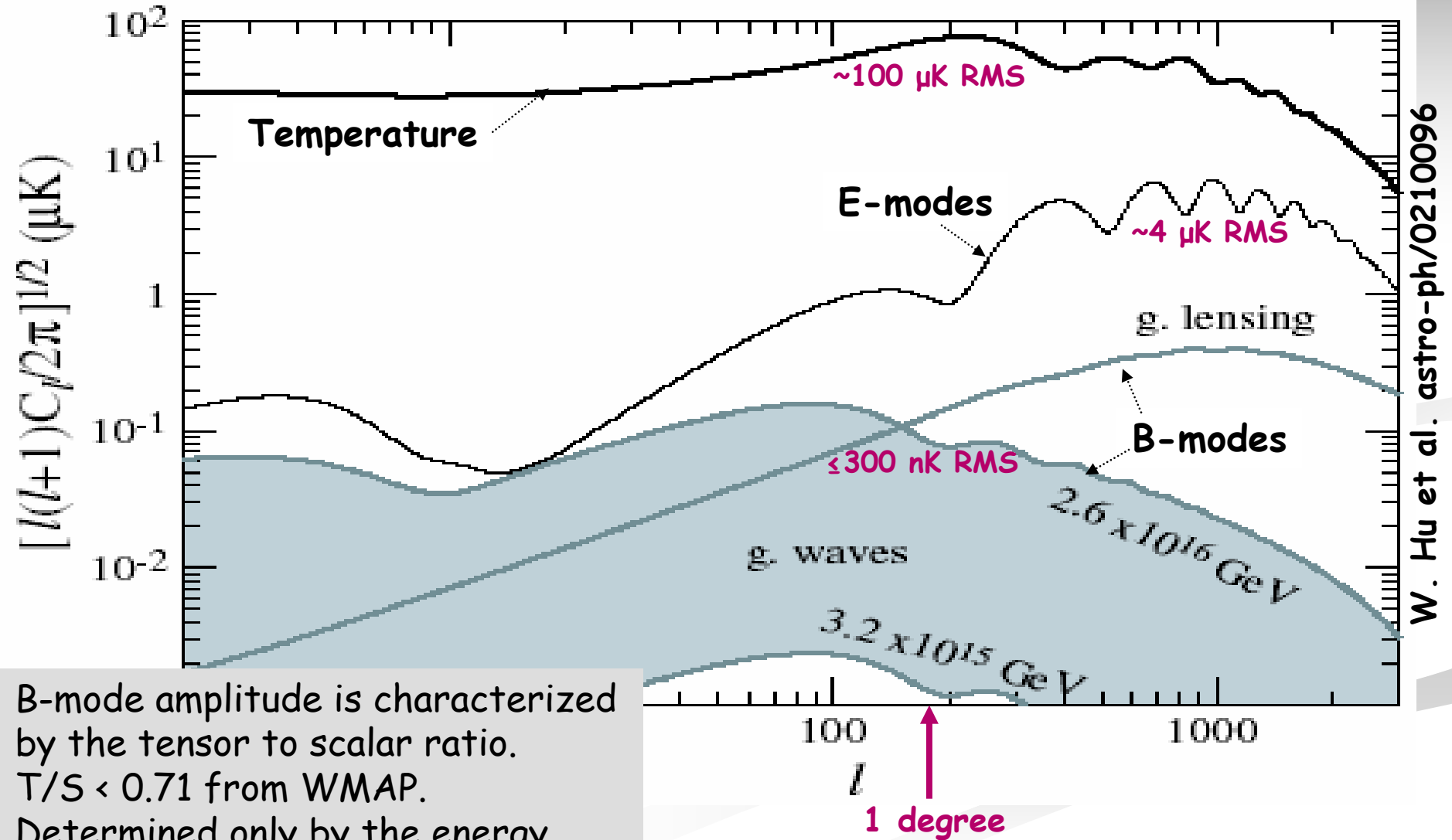
IGB



CMB Polarization

LBNL Teach-In, July 2, 2003 28

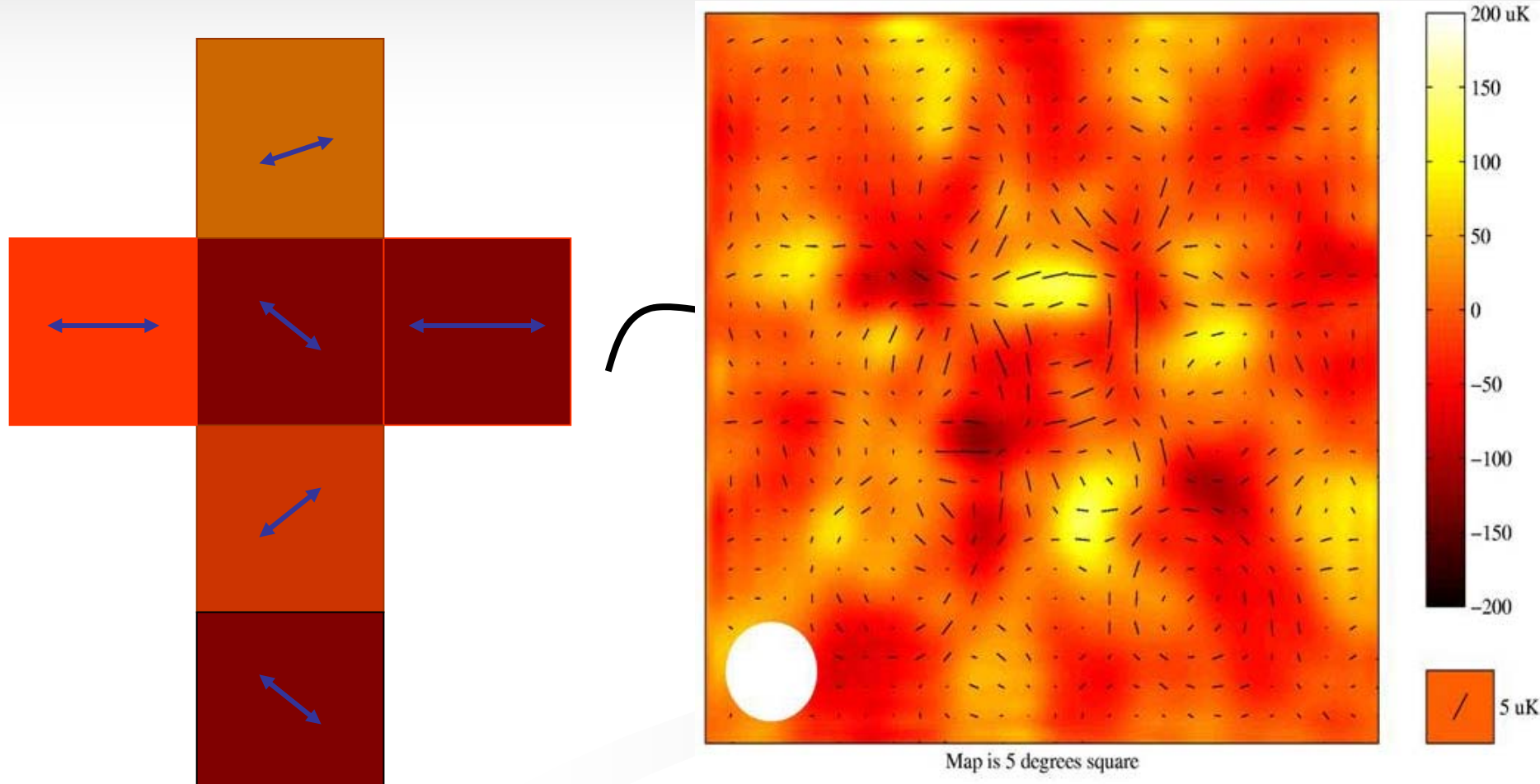
Polarization Power Spectrum



B-mode amplitude is characterized by the tensor to scalar ratio. $T/S < 0.71$ from WMAP. Determined only by the energy scale of inflation.

Polarization Sensitive MAP

- we upgraded our toy experiment to be polarization sensitive.
- For each pixel we measure I,Q,U Stokes parameters, and obtain a map like:



l=2 moment of Polarization Spectrum

$\Delta I = -5\mu K$ $Q = 0.1$ $U = +0.1$ 90, 90	$\Delta I = -5\mu K$ $Q = 0.1$ $U = +0.1$ 90, 90	$\Delta I = -5\mu K$ $Q = 0.1$ $U = +0.1$ 90, 90
$\Delta I = -10\mu K$ $Q = 0.4$ $U = 0$ 0, 0	$\Delta I = +10\mu K$ $Q = 0$ $U = -0.1$ 90, 0	$\Delta I = +10\mu K$ $Q = 0.5$ $U = 0$ 180, 0
$\Delta I = +5\mu K$ $Q = 0$ $U = +0.1$ 90, 270	$\Delta I = +5\mu K$ $Q = 0$ $U = +0.1$ 90, 270	$\Delta I = +5\mu K$ $Q = 0$ $U = +0.1$ 90, 270
$\Delta I = +10\mu K$ $Q = 0$ $U = -0.1$ 90, 180	$\Delta I = +10\mu K$ $Q = 0$ $U = -0.1$ 90, 180	$\Delta I = +10\mu K$ $Q = 0$ $U = -0.1$ 90, 180

← $\Theta(\hat{n}), Q(\hat{n}), U(\hat{n})$

The multipole moments are :

$$\Theta_{lm} = \int d\hat{n} {}_{s=0}Y_{lm}^*(\hat{n}) \Theta(\hat{n})$$

$$E_{lm} \pm iB_{lm} = - \int d\hat{n} {}_{s=2}Y_{lm}^*(\hat{n}) [Q(\hat{n}) \pm iU(\hat{n})]$$

and we have used the spin - weighted spherical harmonics

$$i.e. {}_{s=0}Y_{2,0}(\hat{n}) = \frac{1}{2} \sqrt{\frac{5}{4\pi}} (3 \cos^2 \theta_i - 1) \rightarrow {}_{s=2}Y_{2,0}(\hat{n}) = \frac{3}{4} \sqrt{\frac{5}{6\pi}} (\sin^2 \theta_i)$$

and our two point correlation functions are :

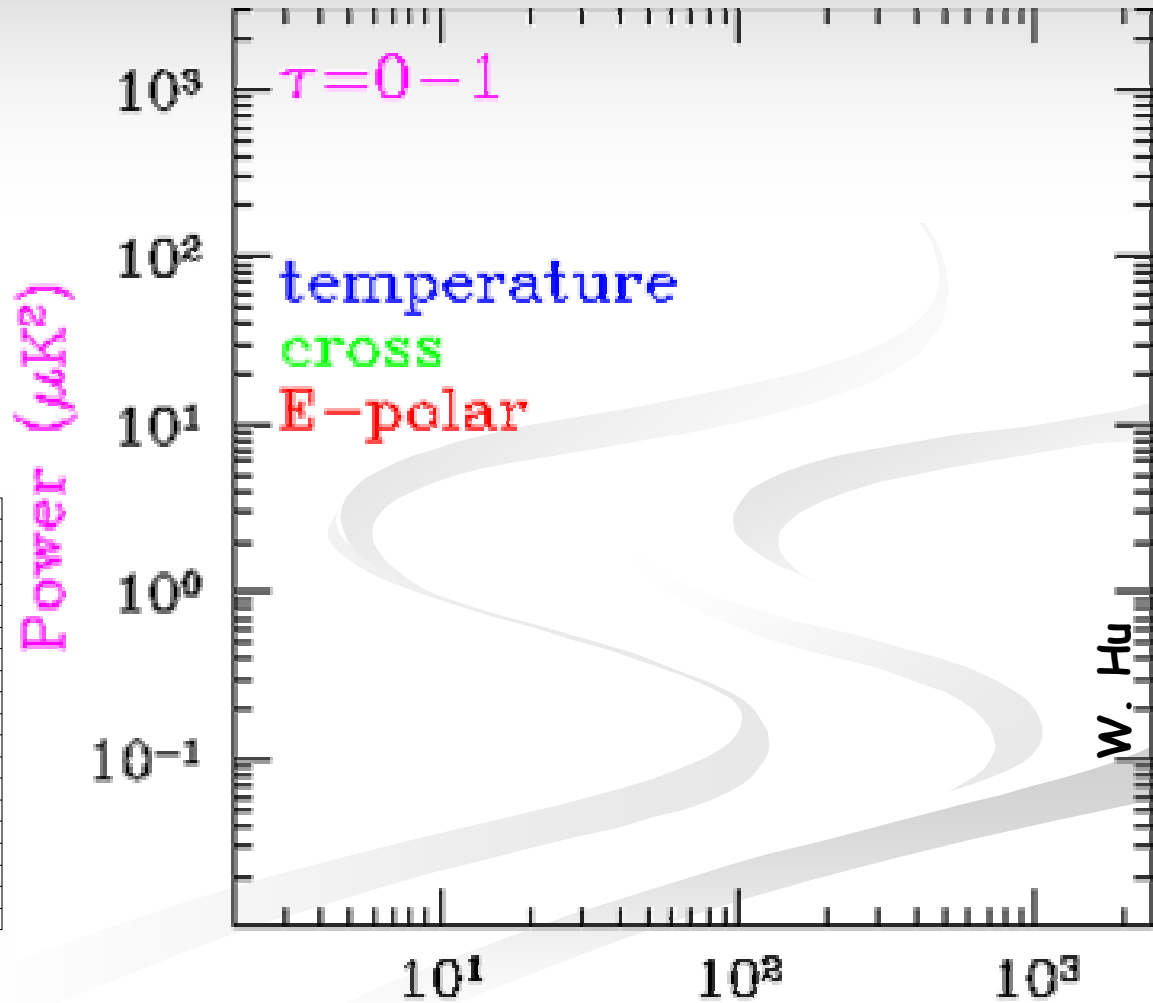
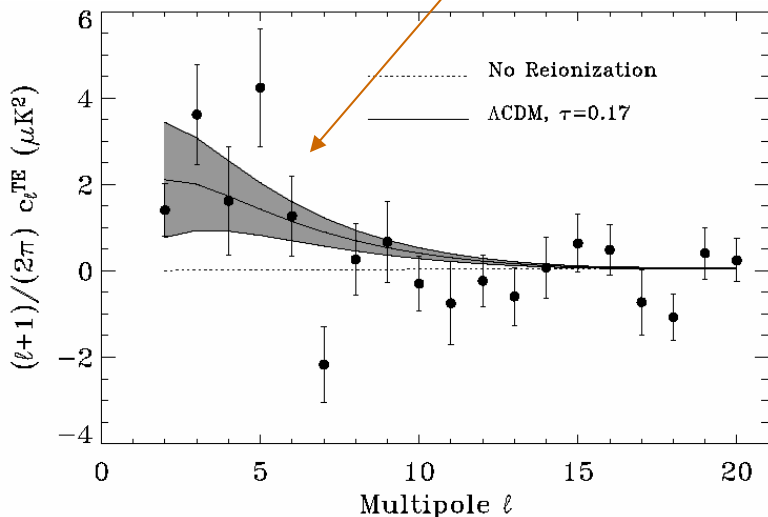
$$\langle X_{lm}^* X'_{l'm'} \rangle = \delta_{ll'} \delta_{mm'} C_l^{XX'} \text{ providing power spectra } C_l^{\Theta\Theta}, C_l^{\Theta E'}, C_l^{EE'}, C_l^{BB'}.$$

($C_l^{BX'}$ with $X' \neq B$ is zero if parity is conserved).

Reionization

- In the temperature $C_l^{\theta\theta}$ there is a degeneracy between reionization, and the overall normalization.
- This degeneracy is broken by the polarization fields.

⊗ Measured by WMAP
 $\tau = 0.017 \pm 0.04$



(a few of) The Polarization Experiments

- DASI (interferometer at the South Pole)
 - ⊗ first observation of E-mode polarization 2002
- Boomerang / MAXIPOL Balloon based telescopes
 - ⊗ E-mode observations
- WMAP NASA Satellite 2001-5
 - ⊗ $C^{\Theta E}_l$ measured, \rightarrow low l bump consistent w/ reionization
 - ⊗ E-mode spectrum expected. B mode reach is small.
- Bicep (South Pole 2004/5)
 - ⊗ E-mode polarization on large angular scales, B-mode discovery potential
- POLARBEAR (California 2005-) 3m ground based telescope
 - ⊗ characterize E-modes, discover B-modes
 - ⊗ facility for testing satellite technology
- Planck ESA/NASA Satellite (2007-)
 - ⊗ characterize E-modes, discovery potential for B-modes
- CMBPOL Einstein Probe Satellite (2013/16)
 - ⊗ definitive measurements of E-mode, significant B-mode reach

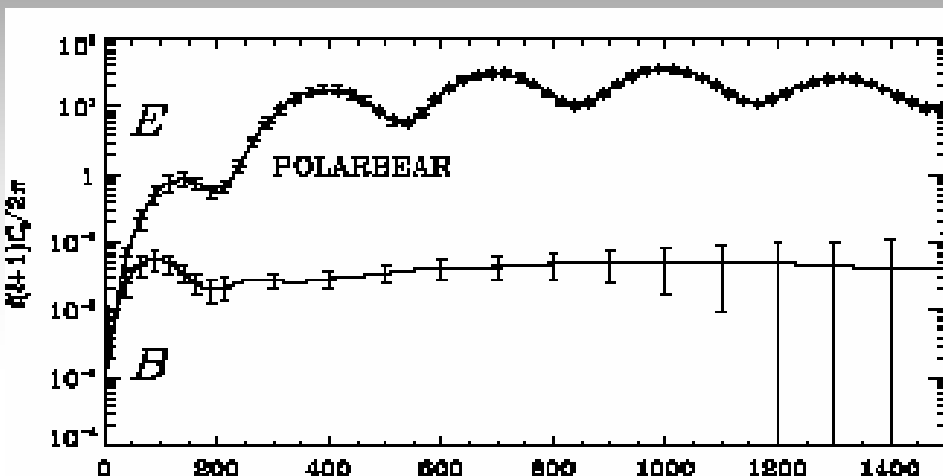
Ground based: POLARBEAR

- LBNL/UCB project (proposed)

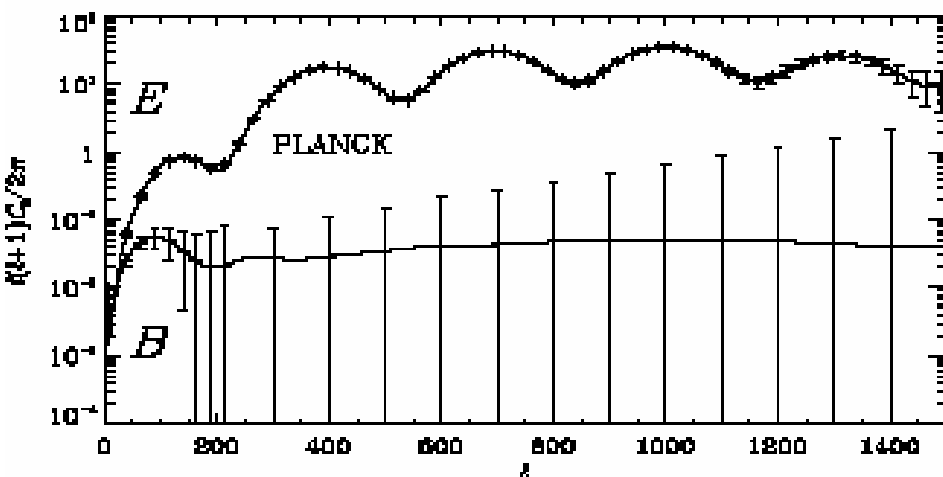


- ❁ first light 2005 from White Mtn with a $\sim 3\text{m}$ primary.
 - atmospheric emission is nearly unpolarized.
 - large sky coverage for primordial gravity waves
 - sufficient resolution to measure and subtract out gravitational lensing signal.
 - staged deployment - 300 elements, upgrade to fMUX 1000s of elements.
 - multi-frequency polarization sensitive antenna coupled to Transition Edge Sensor bolometers
- testing facility for future satellite technologies, systematics, and foreground measurements.

POLARBEAR Performance



3 yrs x 25% duty
 \rightarrow 200 nK/pixel on $10 \times 10 \text{ deg}^2$



3 years observing
 (Assuming a Tensor to Scalar ratio $r=T/S=0.35$)

Sensitivity:

Polarbear I $r=T/S > 0.02$

Polarbear II $r=T/S > 0.0005$

... but understanding systematics
 and foregrounds are the biggest
 challenges now \rightarrow

Polarization Systematics

- instrument polarization $I \rightarrow Q, U$
 - ⊗ smooth term $\delta T_i \rightarrow \delta E_i, \delta B_i$ angle dependent term $\delta T_i \rightarrow \delta E_i, \delta B_i$
 - ⊗ fixed in the instruments coordinate system \rightarrow optimize scan strategy.
- cross polarization $E_x \leftrightarrow E_y \quad Q \leftrightarrow U$
- polarization efficiency (i.e. does polarized light get depolarized or lost) $Q, U \rightarrow I$
- polarization conversion $Q/U \rightarrow \text{signal}$

Our Toolkit:

- ⊗ Multiple levels of differencing:
 - subtraction of orthogonal antennas
 - modulate polarization w/ half wave plate, etc.
 - scan strategy, rotation of sky
- ⊗ High SNR helpful - e.g. compare halves of the data.
- ⊗ low side-lobe response telescope to avoid polarization from ground reflections, etc.

Could measure these effects with a known polarized sources and perfect unpolarized black body.

Polarization Foregrounds

■ Gravitational Lensing

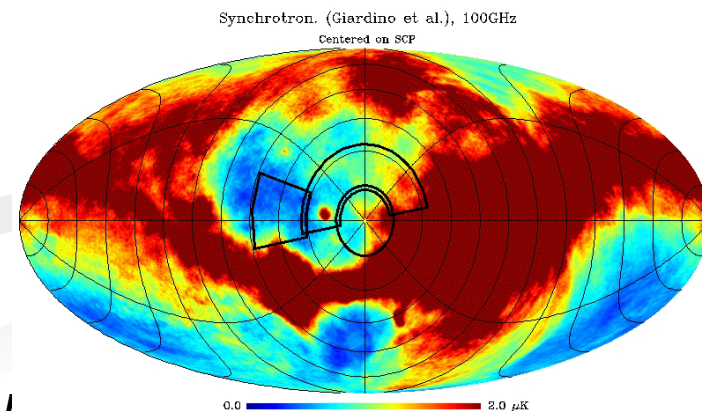
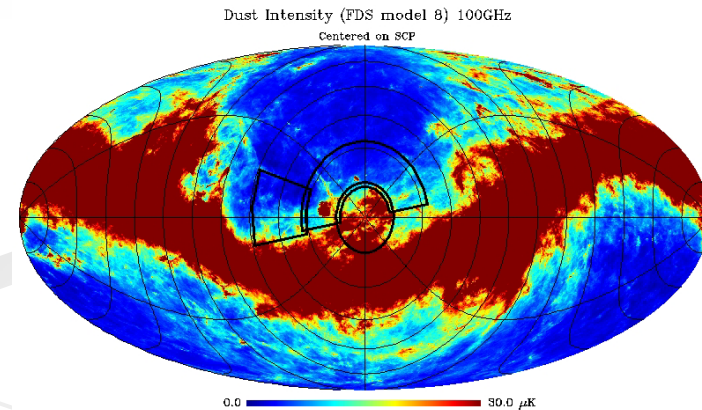
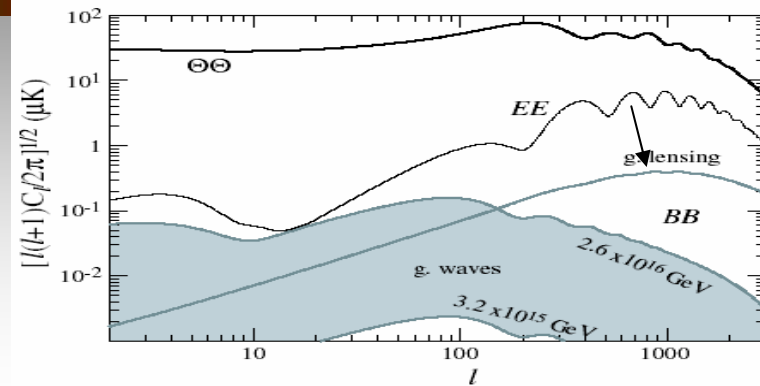
- ✿ converts $E \rightarrow B$ modes
- sufficient angular resolution (high l) is primary tool for subtracting out cosmic shear.

■ Galactic dust

- ✿ intensity is well known
- ✿ polarization very poorly known

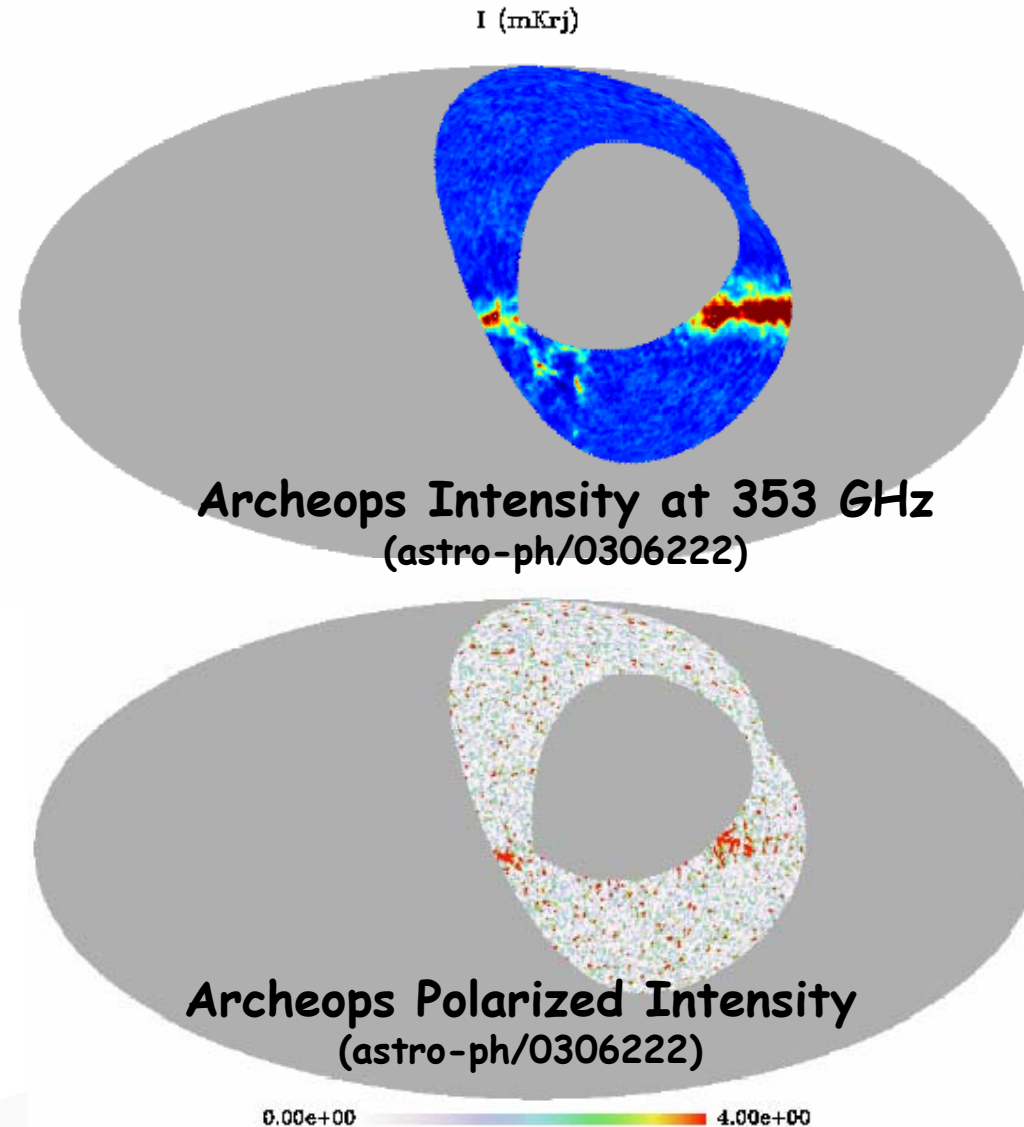
■ Synchrotron Polarized Emission

- ✿ must extrapolate data from 1-3 GHz up to 100 GHz.

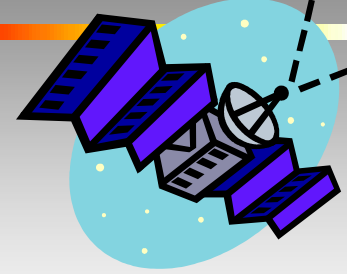


Polarized Galactic Dust Emission

- new results from Archeops (balloon flown over Sweden) provides first measurements of dust-pol on scales $>10'$
 - Typically polarized at 4-5% level, can be as high as 10-20%.
 - These maps look very different in different bands.
- multiband observations are primary tool against dust-pol.

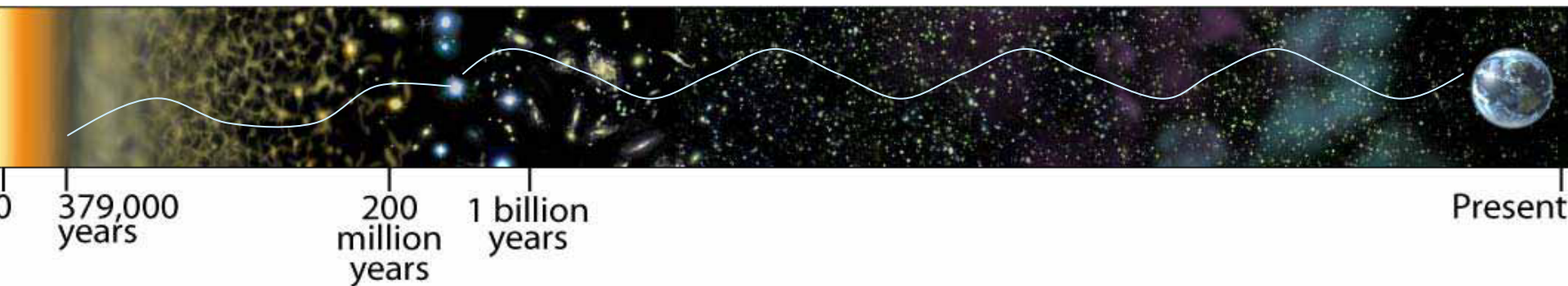


Einstein Probe Satellite: CMBPOL



- Proposals for concept studies submitted.
- 2013-16 launch with science goals
 - ⊗ definitive measurement of E-mode polarization
 - cosmic variance limited all the way out to the damping tail.
 - ⊗ detect B-mode polarization down to $E_{\text{inflation}} < \sim 10^{15} \text{ GeV}$
 - sensitivity limited by residual astrophysical foregrounds
 - ⊗ measure w equation of state with SZ, cosmic shear
 - ⊗ measure sum of neutrino masses with cosmic shear
- expect polarization sensitivity of $1 \mu\text{K s}^{-1}$
- lots of R&D between now and then- upcoming polarization facilities will be test-beds for CMBPOL instrumentation.
 - ⊗

more CMB science...

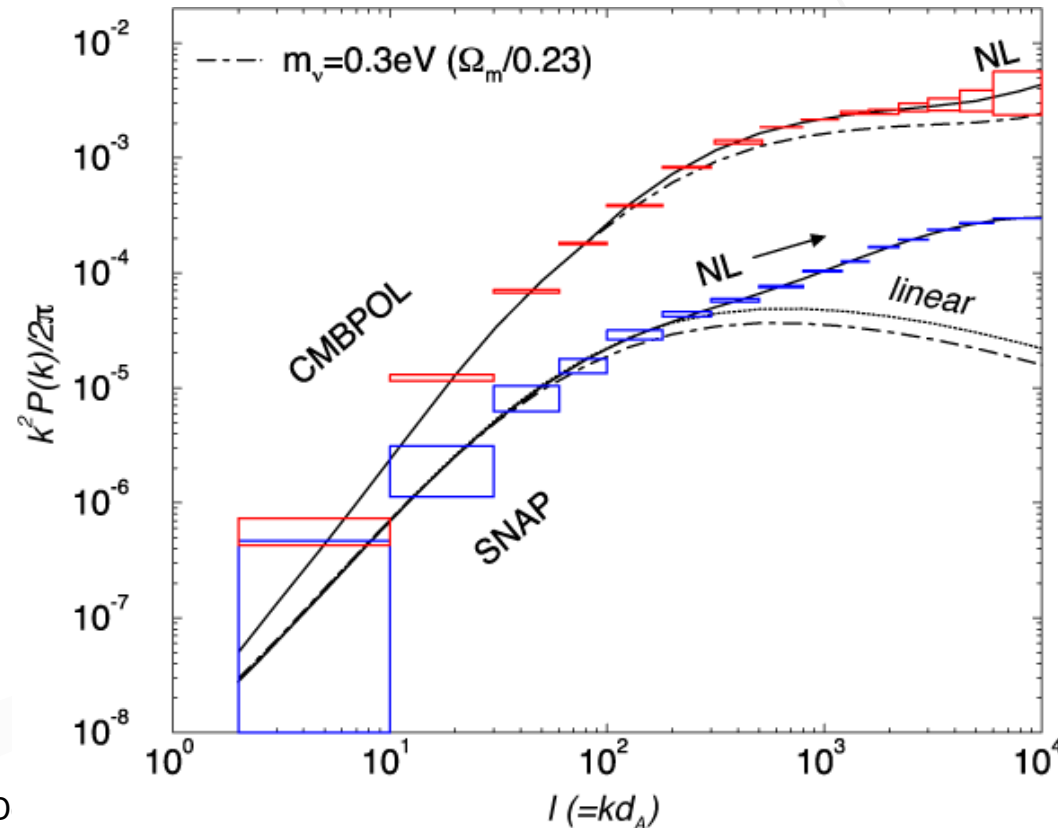


- CMB Secondaries- physics that happens after the primary anisotropy is encoded at the surface of last scattering.
- CMB photons are used to backlight structure in the universe that is newer than $z \sim 1100$.
 - ⊗ Cluster surveys with the Sunyaev Zeldovich effect $\rightarrow \Omega_M, w$
 - ⊗ gravitational lensing
 - measurement of the neutrino mass
 - ⊗ ...

Grav. Lensing of CMB: Neutrino Mass

- power spectrum for cosmic shear of the CMB at small angular scales is sensitive to neutrino masses
 - ⊗ measures the matter content (all matter!, not just luminous matter!) along the light of sight.
 - ⊗ future CMBPOL satellite may have a sensitivity of ~ 0.03 eV.

Neutrino mass effects enter where linear perturbation theory applies... different from SNAP.



Conclusions

- E-mode polarization
 - ⊗ significantly increases info about the cosmo parameters
 - ⊗ removes the degeneracy from reionization
- B-mode polarization
 - ⊗ tool for probing Inflationary Gravity Wave Background
 - ⊗ secondaries provide equally exciting science: matter power spectrum, neutrino masses, cluster surveys
 - ⊗ systematics and foregrounds are important
- new E-mode results expected soon
 - ⊗ Boomerang, Maxipol, WMAP
- first significant B-mode measurements
 - ⊗ will come from the ground, allowing easy modifications, to handle foregrounds and systematics- as they become known
 - ⊗ This knowledge will be assembled to construct a CMB polarization satellite ~ 2015.

Resources for Further Study

- Wayne Hu's Webpage, <http://background.uchicago.edu/~whu/>
- Hu/Dodelson, *Ann. Rev. Astron. Astrophys.* **40** 171 (2002) [astro-ph/0110414](#)
- Hu/White, *A CMB Polarization Primer*, [astro-ph/9706147](#).
- for HEP folks: Rubakov, *Cosmology and Astrophysics*, CERN Summer School, CERN-YR-2002-002.